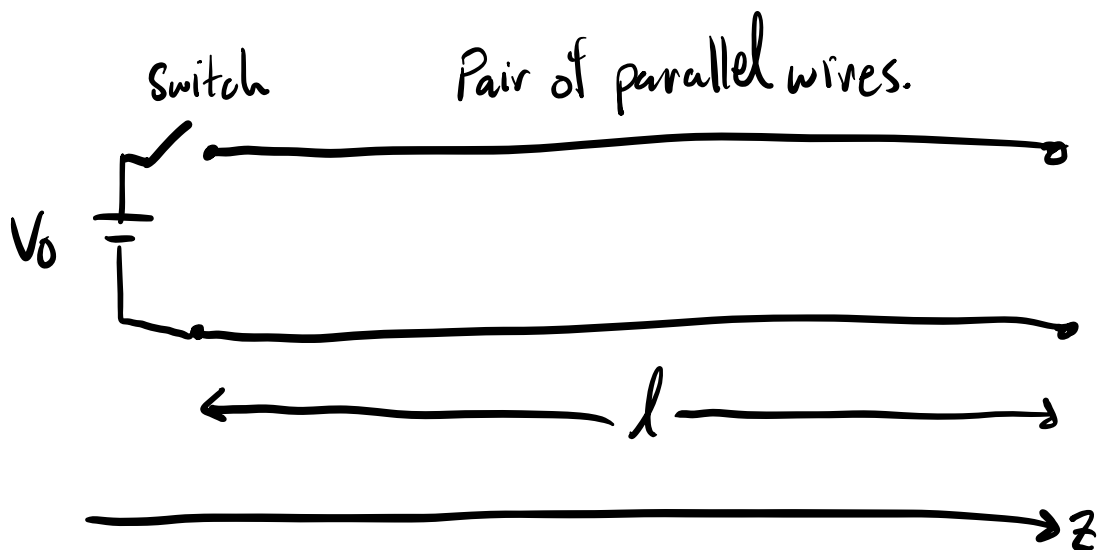
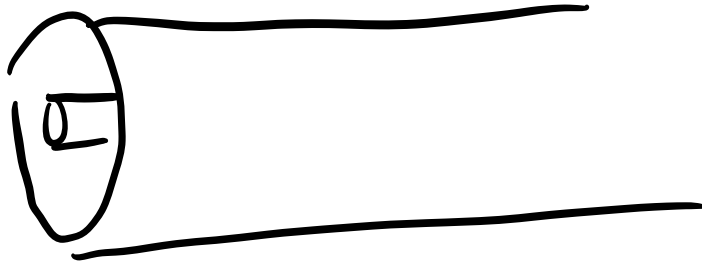


PHYS 331 - September 27, 2022

Transmission Lines.

Coaxial
cable



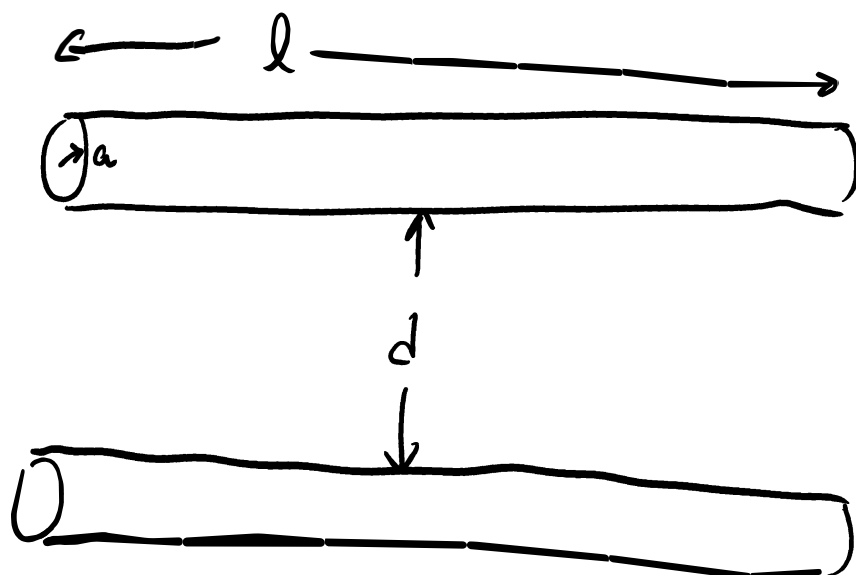
When switch is closed, get a pot. diff. across wires \rightarrow wires store charge \rightarrow capacitance

$$C = \frac{q}{\Delta V}$$

Must be a transient current to transport the charge \rightarrow inductance

$$V_L = L \frac{dI}{dt}$$

$d \gg a$



$$\int \vec{E} \cdot d\vec{A} = \frac{q_{\text{enc}}}{\epsilon_0}$$

$$\Delta V = - \int \vec{E} \cdot d\vec{l}$$

capacitance
per unit length.

$$\frac{C}{l} = C_l = \frac{\pi \epsilon_0}{\ln\left(\frac{d}{a}\right)}$$

$$\int \vec{B} \cdot d\vec{l} = \mu_0 i_{\text{enc}}$$

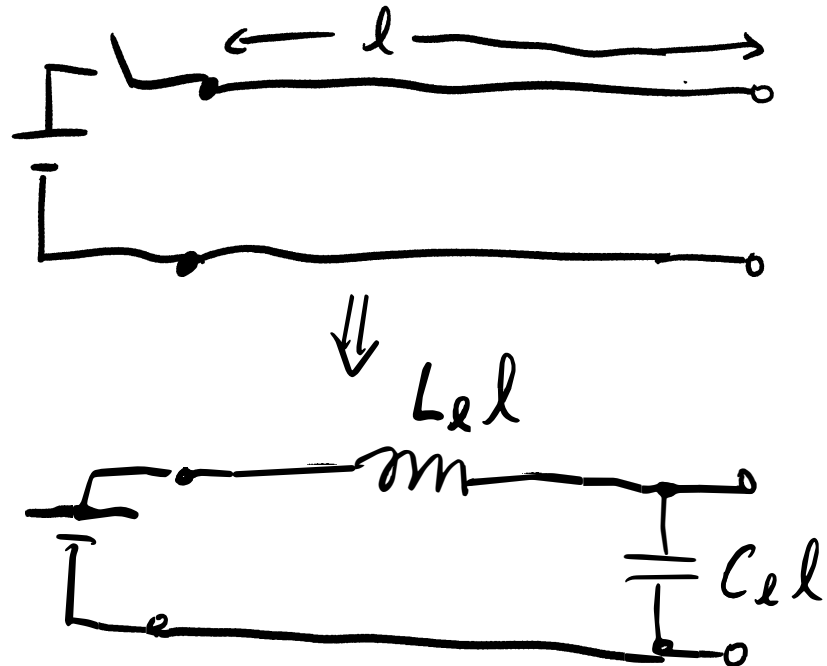
$$\Phi_B = \int \vec{B} \cdot d\vec{A}$$

$$V_L = - \frac{d\Phi_B}{dt} = -L \frac{dI}{dt}$$

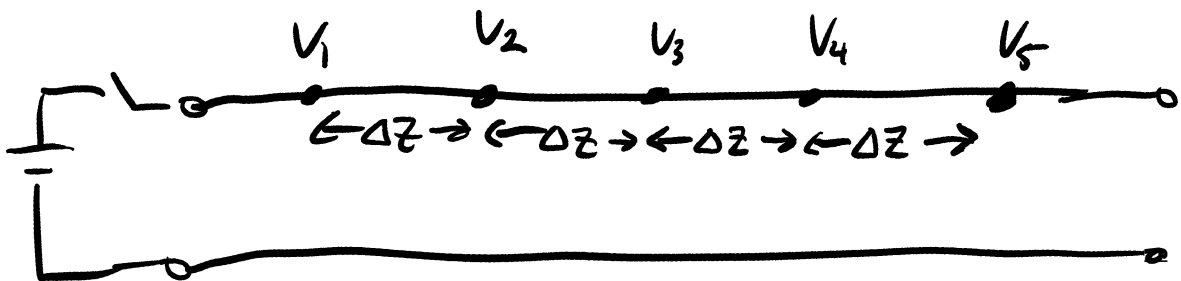
inductance per
unit length

$$\frac{L}{l} = L_l = \frac{\mu_0}{\pi} \ln\left(\frac{d}{a}\right)$$

Model parallel wires as a circuit.



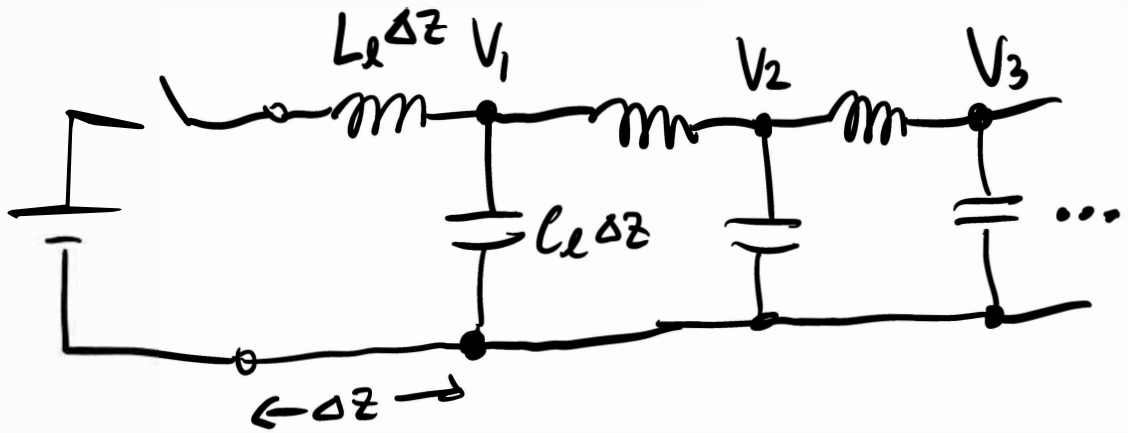
If wires are very long...



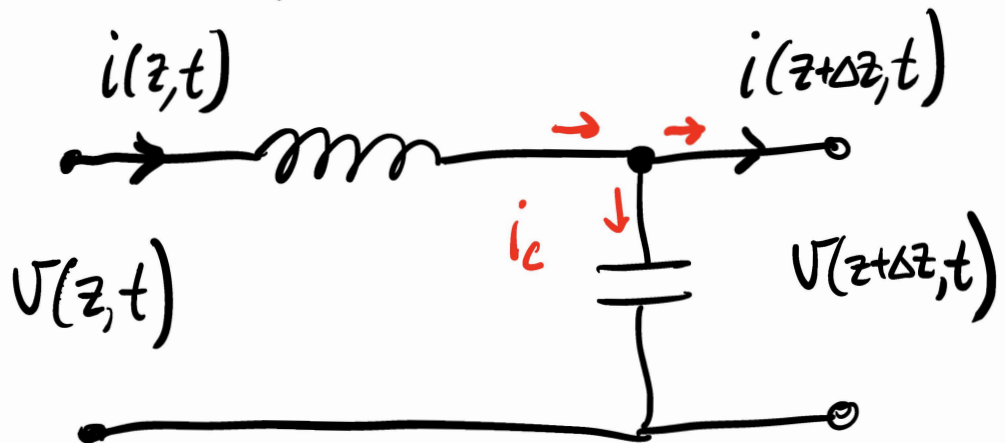
$V_1 > V_2 > V_3 > V_4 > V_5$ a short time after switch is closed.

The simple circuit model fails for long transmission lines.

Model long transmission line as a distributed LC network.



Consider a single branch of this circuit



Kirchhoff loop rule

$$V(z, t) - L_l \Delta z \frac{\partial i(z, t)}{\partial t} = V(z + \Delta z, t)$$

$$\therefore \frac{V(z + \Delta z, t) - V(z, t)}{\Delta z} = -L_l \frac{\partial i(z, t)}{\partial t}$$

In $\Delta z \rightarrow 0$ limit, we have

$$\boxed{\frac{\partial U(z,t)}{\partial z} = -L \frac{\partial i(z,t)}{\partial t}} \quad \textcircled{1}$$

Kirchhoff's Junction Rule:

$$C = \frac{q}{\Delta V_c} \quad \therefore q = C \Delta V_c$$

$$\therefore i_c = C \frac{dV_c}{dt}$$

$$i(z,t) = i_c + i(z+\Delta z, t)$$

$$\text{Now sub } i_c = C \Delta z \frac{\partial U(z+\Delta z, t)}{\partial t}$$

$$\therefore i(z,t) = C \Delta z \frac{\partial U(z+\Delta z, t)}{\partial t} + i(z+\Delta z, t)$$

$$\therefore \frac{i(z+\Delta z, t) - i(z,t)}{\Delta z} = -C \frac{\partial U(z+\Delta z, t)}{\partial t}$$

In limit $\Delta z \rightarrow 0$

$$\boxed{\frac{\partial i(z,t)}{\partial z} = -L_e \frac{\partial v(z,t)}{\partial t}} \quad (2)$$

Assume that v & i are harmonic

$$\rightarrow v(z,t) = V(z)e^{j\omega t}$$

$$\rightarrow i(z,t) = I(z)e^{j\omega t}$$

From (1)

$$\frac{\partial}{\partial z} (V(z)e^{j\omega t}) = -L_e \frac{\partial}{\partial t} (I(z)e^{j\omega t})$$

$$\frac{\partial V(z)}{\partial z} \cancel{e^{j\omega t}} = -j\omega L_e I(z) \cancel{e^{j\omega t}}$$

$$\boxed{\frac{\partial V}{\partial z} = -j\omega L_e I} \Rightarrow \frac{\partial^2 V}{\partial z^2} = -j\omega L_e \frac{\partial I}{\partial z}$$

From (2)

$$\frac{\partial}{\partial z} (I(z)e^{j\omega t}) = -C_e \frac{\partial}{\partial t} (V(z)e^{j\omega t})$$

$$\frac{\partial I}{\partial z} \cancel{e^{j\omega t}} = -j\omega C_e V(z) \cancel{e^{j\omega t}}$$

$$\boxed{\frac{\partial I}{\partial z} = -j\omega C_e V} \Rightarrow \frac{\partial^2 I}{\partial z^2} = -j\omega C_e \frac{\partial V}{\partial z}$$

$$\therefore \frac{\partial^2 V}{\partial z^2} = -j\omega L_e (-j\omega C_e V)$$

$$\boxed{\therefore \frac{\partial^2 V}{\partial z^2} = -\omega^2 L_e C_e V}$$

The wave Eq'ns. ↓

Likewise:

$$\frac{\partial^2 I}{\partial z^2} = -j\omega C_e (-j\omega L_e I)$$

$$\frac{\partial^2 I}{\partial z^2} = -\omega^2 L_e C_e I$$

General sol'n to voltage eq'n:

$$V(z) = V_+ e^{-j\beta z} + V_- e^{j\beta z}$$

check:

$$\frac{\partial V}{\partial z} = -j\beta V_+ e^{-j\beta z} + j\beta V_- e^{j\beta z}$$

$$\begin{aligned} \frac{\partial^2 V}{\partial z^2} &= -\beta^2 V_+ e^{-j\beta z} - \beta^2 V_- e^{j\beta z} \\ &= -\beta^2 \underbrace{(V_+ e^{-j\beta z} + V_- e^{j\beta z})}_{V(z)} = -\beta^2 V(z) \end{aligned}$$

Assumed sol'n for $V(z)$ is valid provided

$$\beta^2 = \omega^2 L_e C_e$$

$$\beta = \omega \sqrt{L_e C_e}$$

Full time-dependent voltage is:

$$V(z, t) = V(z) e^{j\omega t}$$

$$= (V_+ e^{-j\beta z} + V_- e^{j\beta z}) e^{j\omega t}$$

$$= V_+ e^{j(\omega t - \beta z)} + V_- e^{j(\omega t + \beta z)}$$